

# Large-Area Synthesis of Continuous and Uniform MoS<sub>2</sub> Monolayer Films on Graphene

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Heterostructures composed of multiple layers of different atomically thin materials are of interest due to their unique properties and potential for new device functionality. MoS<sub>2</sub>-graphene heterostructures have shown promise as photodetectors and vertical tunnel transistors. However, progress is limited by the typically micrometer-scale devices and by the multiple alignments required for fabrication when utilizing mechanically exfoliated material. Here, the synthesis of large-area, continuous, and uniform MoS<sub>2</sub> monolayers directly on graphene by chemical vapor deposition is reported, resulting in heterostructure samples on the centimeter scale with the possibility for even larger lateral dimensions. Atomic force microscopy, photoluminescence, X-ray photoelectron, and Raman spectroscopies demonstrate uniform single-layer growth of stoichiometric MoS<sub>2</sub>. The ability to reproducibly generate large-area heterostructures is highly advantageous for both fundamental investigations and technological applications.

## 1. Introduction

Graphene has received much attention due to its intrinsic two-dimensional (2D) nature, high electronic mobility, linear band dispersion and exceptional crystalline quality.<sup>[1]</sup> In addition to fundamental scientific interest, it also promises to have significant technological impact on a variety of applications, including flexible electronics,<sup>[2]</sup> hydrogen storage,<sup>[3]</sup> gas sensing,<sup>[4,5]</sup> use as a conductive back plane,<sup>[6]</sup> and protective coatings.<sup>[7]</sup>

This diverse activity has stimulated interest into other 2D materials such as the transition metal dichalcogenides (e.g., MoS<sub>2</sub>, MoSe<sub>2</sub>, WS<sub>2</sub>, and WSe<sub>2</sub>).<sup>[8,9]</sup> These materials share the van der Waals-bonded layered structure of graphite, enabling isolation of single to few monolayer samples for analysis by mechanical<sup>[10]</sup> or chemical exfoliation.<sup>[11]</sup> Monolayer samples thus prepared exhibit properties that are strikingly different from the bulk and complement those of graphene, providing a suite of 2D materials suitable for a much broader range of

applications. For instance, although bulk MoS<sub>2</sub> is an indirect gap semiconductor ( $E_g \approx 1.3$  eV), in monolayer format it abruptly transitions to a direct gap semiconductor with strong excitonic luminescence at  $\approx 1.8$  eV.<sup>[12,13]</sup> Monolayer MoS<sub>2</sub> exhibits high optical responsivity and a quantum yield of  $4 \times 10^{-3}$ ,<sup>[12]</sup> in direct contrast to graphene, which displays inefficient light emission and quantum efficiency below  $10^{-12}$ .<sup>[14]</sup> Transistors fabricated from MoS<sub>2</sub> monolayers demonstrate room temperature current on/off ratios in excess of  $1 \times 10^8$  and ultralow standby power dissipation,<sup>[15]</sup> as well as highly selective chemical vapor sensors.<sup>[16]</sup> These properties suggest MoS<sub>2</sub> may be a promising material for integration into CMOS electronics and a valuable complement to graphene.<sup>[17–20]</sup> By combining graphene

with MoS<sub>2</sub> it may be possible to utilize beneficial properties of each material.

A logical progression from the study of individual 2D materials leads to the combination of multiple films to form heterostructures. Very recently, heterostructure devices were successfully fabricated and characterized,<sup>[21–24]</sup> exhibiting unique and interesting properties. The MoS<sub>2</sub>-graphene combination holds promise for use as a photodetector, as high photocurrents and persistent photoconductivity were demonstrated.<sup>[23]</sup> MoS<sub>2</sub> can also serve as a tunnel barrier for current injection into graphene in a field effect tunneling transistors,<sup>[22]</sup> a difficult task to accomplish on graphene due to its high surface diffusion and low chemical reactivity.

Typical heterostructure devices are produced via mechanical exfoliation and multiple stackings of the various 2D components,<sup>[22–24]</sup> limiting fabrication to a few micrometer-scale devices. Direct large-area growth of the desired 2D heterostructure is critical to enable large-scale device fabrication for technological applications. Chemical vapor deposition (CVD) is a successful means to produce large-area graphene,<sup>[2,25,26]</sup> with current technologies growing continuous films hundreds of meters in length using a roll-to-roll process.<sup>[27]</sup> Additionally, continuous films of MoS<sub>2</sub> were recently grown on insulating substrates.<sup>[28–30]</sup> While the synthesis of individual 2D materials was successfully demonstrated, the growth of heterostructures composed of multiple atomically thin 2D materials remains a key goal and challenge. Several recent reports provided important contributions to the nascent field of 2D heterostructure

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synthesis. Nano-scale flakes of MoS<sub>2</sub> on graphene were achieved<sup>[31,32]</sup> as well as continuous MoS<sub>2</sub> films of varying thickness on epitaxial graphene,<sup>[33]</sup> and seeding promoted monolayer MoS<sub>2</sub> on mechanically exfoliated graphene.<sup>[34]</sup> Although progress is rapid, difficulties remain in the synthesis of large-area heterostructures with uniformity and thickness control.

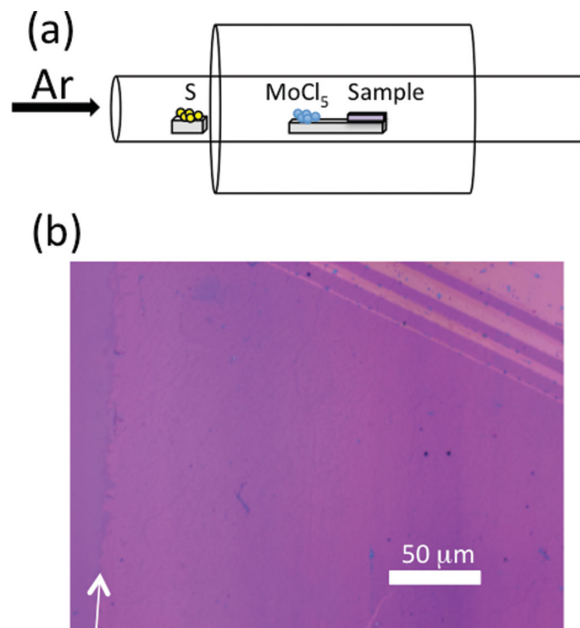
## 2. Results and Discussion

In this paper, we report the growth by CVD of continuous and uniform MoS<sub>2</sub> single layer films on large-area graphene. The heterostructure films are characterized via atomic force microscopy (AFM), Raman spectroscopy, X-ray photoelectron spectroscopy (XPS), and photoluminescence (PL) spectroscopy. Direct large-area production of van der Waals heterostructures will make fundamental investigations of these systems more accessible to the scientific community, and is an important step towards incorporation of such devices into future technology.

A two-step CVD process is implemented to obtain heterostructure samples consisting of MoS<sub>2</sub> on graphene, with a supporting substrate of SiO<sub>2</sub>/Si(001). In the first step, large-area, high quality graphene is grown by CVD on copper foil and transferred onto an SiO<sub>2</sub>(275nm)/Si wafer following the procedure of Li et al.<sup>[25]</sup> The samples are annealed in forming gas to clean the graphene surface and improve adhesion. The graphene/SiO<sub>2</sub>/Si is then introduced to a second tube furnace for the CVD growth of MoS<sub>2</sub>. Solid sources of MoCl<sub>5</sub> and sulfur are used as precursors for the MoS<sub>2</sub> growth.<sup>[28]</sup> Both the substrate and MoCl<sub>5</sub> precursor are placed on a quartz platform at the center of a 2" diameter tube furnace, with a separation of 4–8 cm between MoCl<sub>5</sub> precursor and substrate, as shown in the schematic of Figure 1a. To achieve monolayer films of MoS<sub>2</sub> on graphene, 20 mg of MoCl<sub>5</sub> is used, with 1200 mg sulfur positioned at the upstream end of the furnace. Thicker films can be controllably synthesized by increasing the amount of MoCl<sub>5</sub>. Two (three) layers are achieved using 35 mg (50 mg) of precursor. As the temperature is increased, the pressure is maintained at 2 Torr under a 50 sccm flow of Argon. Upon reaching 850 °C, the growth temperature is maintained for 10 min, and the sample is then allowed to cool undisturbed under continuous Ar flow.

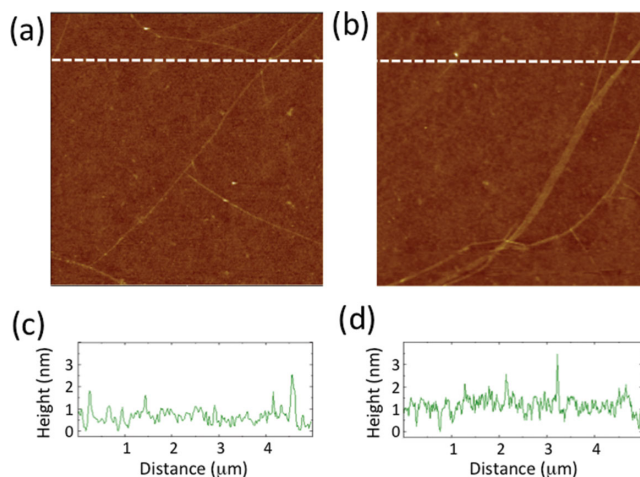
Figure 1b shows an optical image of the resulting growth on graphene. The edge of the underlying graphene is identified via an arrow as a guide to the eye. The MoS<sub>2</sub> uniformly covers the graphene film and exposed SiO<sub>2</sub> areas, as evident in the optical image. The MoS<sub>2</sub>-graphene film covers the area to the right of the arrow, while the MoS<sub>2</sub> on exposed SiO<sub>2</sub> is to the left of the arrow. An intentional scratch following MoS<sub>2</sub> growth highlights the contrast between the MoS<sub>2</sub>-graphene film and bare SiO<sub>2</sub> (upper right corner, Figure 1b). The MoS<sub>2</sub> forms a continuous layer on the entirety of the substrate, indicating that heterostructure growth could be scaled to larger dimensions.

AFM is utilized to investigate the topography of the as-grown MoS<sub>2</sub>/graphene/SiO<sub>2</sub>/Si heterostructures. Figure 2a displays a 5 μm × 5 μm image of the MoS<sub>2</sub>-graphene film, while Figure 2b shows an image of the graphene film alone. The ripples and wrinkles present in large-area graphene (Figure 2b) are also apparent in the MoS<sub>2</sub> grown on graphene, indicating that the

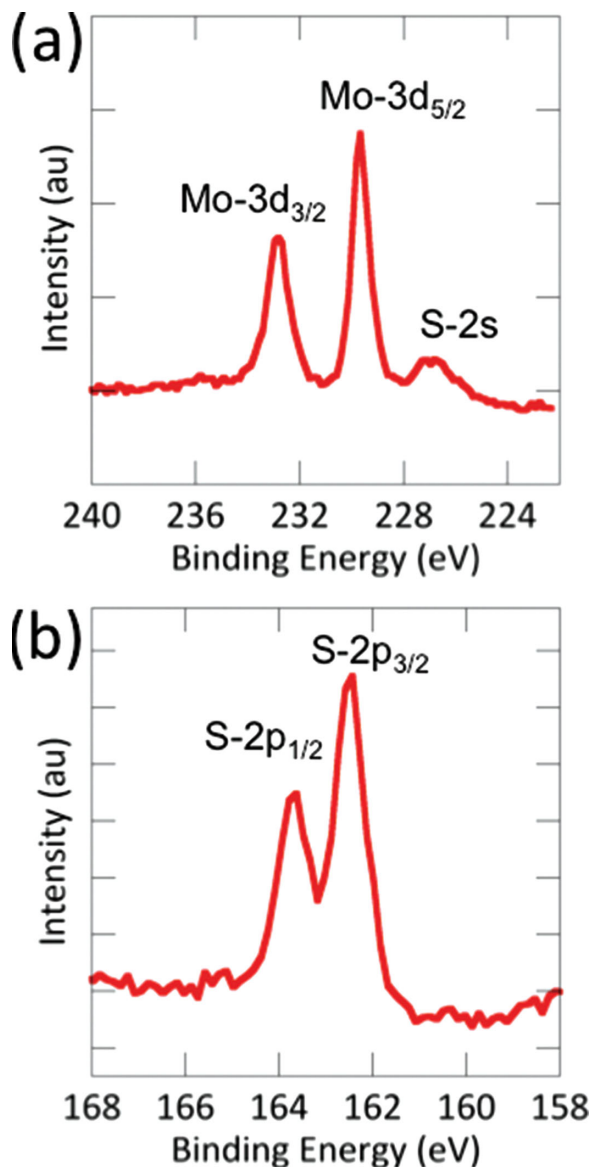


**Figure 1.** a) Schematic of the furnace set-up for MoS<sub>2</sub> growth on graphene. b) Optical image following the growth procedure. The resulting MoS<sub>2</sub> is continuous and uniform on graphene (to the right of the arrow) as well as on SiO<sub>2</sub> (to the left of the arrow). The heterostructure is intentionally scratched in the upper right corner to indicate contrast with bare SiO<sub>2</sub>.

MoS<sub>2</sub> growth process results in a uniform film that follows the contours of the substrate. The height profiles along the surface of MoS<sub>2</sub>-graphene (Figure 2c) and bare graphene (Figure 2d) are nearly analogous. We observe small clusters of thicker materials attributed to residues that have not been fully removed during the graphene cleaning procedure, rather than non-uniform MoS<sub>2</sub> growth, as similar clusters are apparent on bare graphene (Figure 2b). Both MoS<sub>2</sub>-graphene and bare graphene exhibit an rms roughness less than 0.5 nm. We observe no variations in the number of layers, as confirmed by photoluminescence and



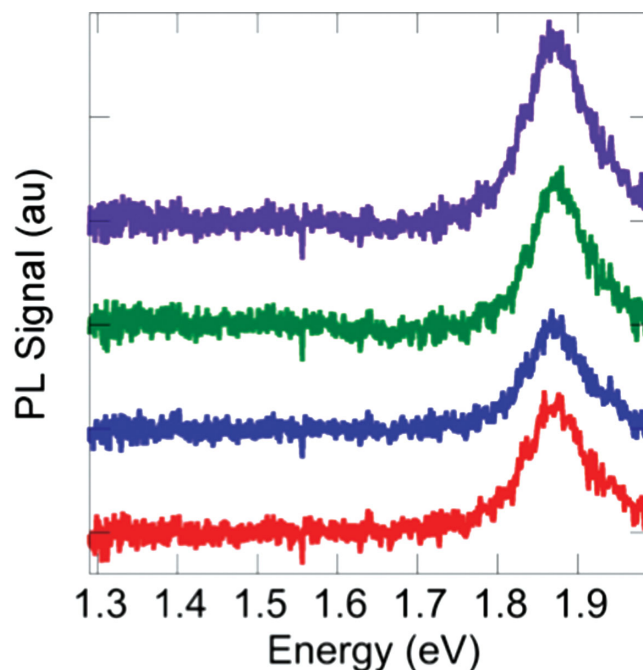
**Figure 2.** AFM images of a 5 μm × 5 μm area. a) The image of the MoS<sub>2</sub>-graphene heterostructure and b) graphene substrate for comparison. c,d) Height profiles along the dotted line in (a) and (b), respectively. The rms roughness is below 0.5 nm before and after MoS<sub>2</sub> growth.



**Figure 3.** Core binding energies determined using XPS spectra. a) Mo 3d peaks are present at 232.6 eV and 229.5 eV, with S 2s at 226 eV. b) S peaks are present at 163.5 eV and 162.2 eV. XPS indicate stoichiometric MoS<sub>2</sub> synthesis on graphene.

Raman spectroscopy (discussed below), nor any indication that ripples or edges of graphene serve as nucleation sites for the MoS<sub>2</sub> growth.

X-ray photoelectron spectroscopy is used to probe the chemical composition of the as-grown films (which are exposed to atmosphere before loading into the XPS system). Spectra are obtained with monochromated Al K $\alpha$  radiation using a 400  $\mu$ m aperture. **Figure 3a** displays Mo 3d peaks at 232.6 eV and 229.5 eV, corresponding to the 3d<sub>5/2</sub> and 3d<sub>3/2</sub> doublet, as well as the sulfur 2s peak at 226 eV. The peaks at 163.5 eV and 162.2 eV (**Figure 3b**) correspond to the S p<sub>1/2</sub> and S p<sub>3/2</sub> orbital. The observed binding energies are consistent with previously reported values for MoS<sub>2</sub><sup>[29,35]</sup> and confirm the expected charge states of S<sup>2-</sup> and Mo<sup>4+</sup> in the MoS<sub>2</sub>. The Mo 3d peaks show no



**Figure 4.** Photoluminescence spectra from the heterostructure samples. The four displayed spectra (offset for clarity) span a 2 mm distance and are representative of the 2.5 mm  $\times$  2.5 mm area investigated. The single emission peak at 1.88 eV and absence of lower emission peaks confirms atomically thin MoS<sub>2</sub> grown on graphene.

indication of additional peaks at lower binding energy (indicative of metallic Mo) or at higher binding energy (indicative of MoO<sub>2</sub> or MoO<sub>3</sub>) attesting to the quality of the growth. Furthermore, XPS spectra measure an S:Mo ratio of 2, confirming stoichiometric MoS<sub>2</sub> synthesis on graphene. Instrumentation limitations and fitting procedures lead to a conservative error of less than 5% in the stoichiometry value.

The MoS<sub>2</sub> on graphene samples are further investigated with photoluminescence spectroscopy measurements. PL is a useful technique for readily distinguishing single and multi-layer MoS<sub>2</sub>. The PL emission of monolayer MoS<sub>2</sub> at room temperature is characterized by a single peak, typically between 1.82 and 1.89 eV. The exact energy is sample dependent and is influenced by both the environment<sup>[36]</sup> and the synthesis method.<sup>[28,29]</sup> Emission from two or more layers of MoS<sub>2</sub> shifts the peak to lower energy, and the intensity is reduced significantly. For multiple layers, the indirect gap transition starts to develop around 1.3 eV.<sup>[12,13]</sup> PL spectra are measured at room temperature over a 2.5 mm  $\times$  2.5 mm area at discrete steps of 0.5 mm. **Figure 4** displays four consecutive measurements, offset for clarity, spanning a lateral distance of 2 mm. All spectra exhibit a strong single-energy emission peak at 1.88 eV with no indication of lower energy peaks and linewidth of  $\approx$ 120 meV. This is consistent with monolayer MoS<sub>2</sub> grown with CVD and further confirms the uniformity of the films.

Raman spectroscopy is a useful tool to investigate properties of the graphene underlayer in addition to MoS<sub>2</sub>. The two characteristic peaks in MoS<sub>2</sub> are the E<sub>2g</sub><sup>1</sup> and A<sub>1g</sub> modes, corresponding to the in-plane and out-of-plane vibrations of the Mo-S bond, respectively. Bulk MoS<sub>2</sub> material exhibits peaks at  $\approx$ 382 cm<sup>-1</sup> and



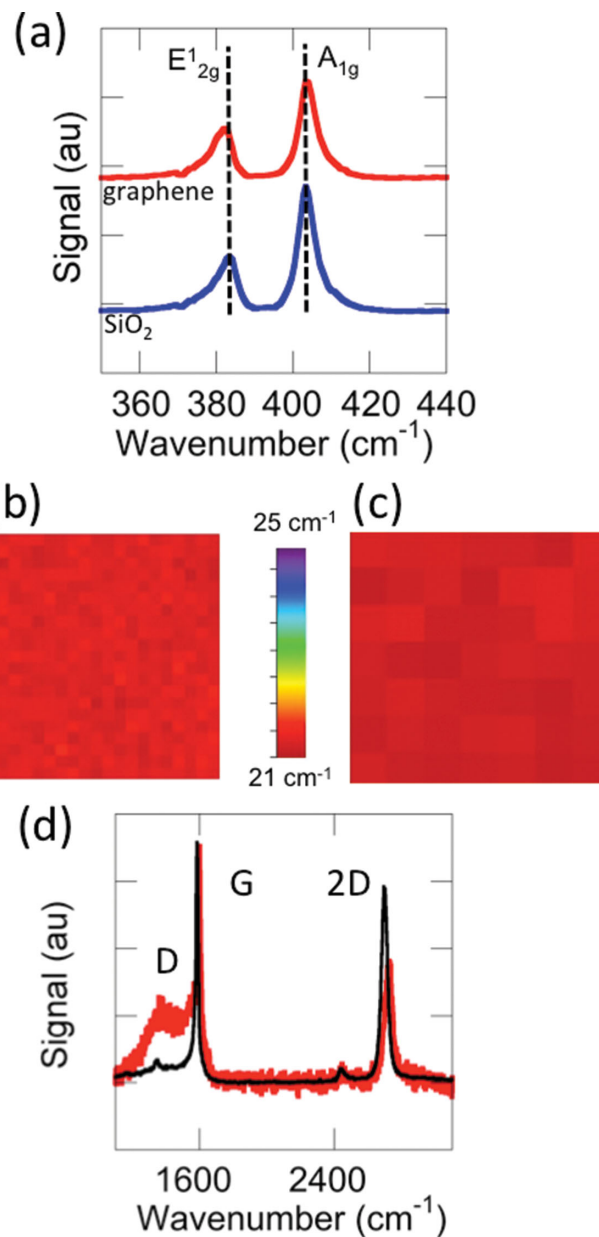
$\approx 407\text{ cm}^{-1}$  with a peak separation ( $\Delta k$ ) of  $25\text{ cm}^{-1}$ .<sup>[37]</sup> As  $\text{MoS}_2$  films approach monolayer thickness,  $\Delta k$  is found to decrease, and the separation between peaks can be used to identify the number of  $\text{MoS}_2$  layers. Experimentally measured  $\Delta k$  for monolayer  $\text{MoS}_2$  on an  $\text{SiO}_2$  substrate are typically in the range of  $18\text{--}20\text{ cm}^{-1}$ , with bilayer increasing to  $\approx 22\text{ cm}^{-1}$ , and trilayer to  $\approx 23\text{ cm}^{-1}$ .<sup>[30,37,38]</sup>

Raman spectra are measured at room temperature using a  $488\text{ nm}$  laser with a spot size  $\approx 0.4\text{ }\mu\text{m}$  and power below  $1\text{ mW}$ . A typical Raman spectrum is displayed as the upper trace in Figure 5a, and exhibits a  $\Delta k$  of  $21.3\text{ cm}^{-1}$ .  $\text{MoS}_2$  grown concurrently on bare  $\text{SiO}_2$  witness substrate exhibits  $\Delta k \approx 20\text{ cm}^{-1}$  as shown in the lower trace, typical of single layer  $\text{MoS}_2$  on an  $\text{SiO}_2$  substrate. The larger peak separation observed in our heterostructure samples is a consequence of the graphene substrate, and has been observed in other experimental works investigating exfoliated  $\text{MoS}_2$ -graphene samples.<sup>[39,40]</sup> The peak separation of  $\approx 21\text{ cm}^{-1}$  observed for our CVD-synthesized  $\text{MoS}_2$ -graphene heterostructures is thus indicative of monolayer  $\text{MoS}_2$  growth, as corroborated by the PL measurements discussed above.

To examine the uniformity across the heterostructure sample on the typical device scale, spectra are acquired at  $0.5\text{ }\mu\text{m}$  spacing across a  $10\text{ }\mu\text{m} \times 10\text{ }\mu\text{m}$  area (Figure 5b). The positions of both  $E_{2g}^1$  and  $A_{1g}$  are constant across the sample surface, with variation in the peak separation  $\leq 0.5\text{ cm}^{-1}$ . Furthermore, the uniformity persists over much larger areas, as can be seen in the  $3\text{ mm} \times 3\text{ mm}$  area displayed in Figure 5c.

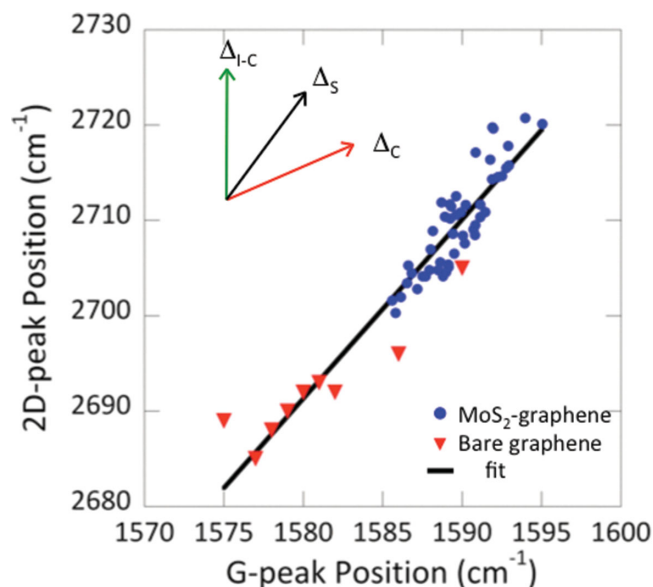
Raman also provides information on how the overgrowth of  $\text{MoS}_2$  impacts the underlying graphene. The graphene peaks of interest are the G, 2D and D peaks as labeled in Figure 5d, with the D peak indicating the presence of defects such as edge states or vacancies. Prior to  $\text{MoS}_2$  growth, we observe all three peaks (black line in Figure 5d), with a very low D peak intensity. Peak positions and the Gaussian lineshape of the 2D peak confirm monolayer graphene.<sup>[41]</sup> Following the  $\text{MoS}_2$  growth (red line in Figure 5d) a small increase in the D/G ratio is observed as well as a shift towards higher wavenumber for both 2D and G peaks. The  $\text{MoS}_2$  film synthesis on top slightly attenuates the Raman signal, resulting in a lower signal to noise ratio in the graphene spectra. Control samples have ruled out elevated temperature as the source of defects, indicating the interaction with the chemical precursors results in defect formation, either as vacancies or  $\text{sp}^3$  bonding. Sulfur has been shown to interact with graphene at elevated temperatures,<sup>[42]</sup> and is the likely source of the enhanced D peak following  $\text{MoS}_2$  growth.

Raman spectroscopy can also indicate charge doping, strain, or a modification of graphene's electronic band structure due to van der Waals forces which we refer to as interlayer coupling. A modulation of electronic band structure near the Dirac point modifies the Fermi velocity,  $v_F$ , of graphene and is apparent in the Raman spectra. Both charge doping and strain produce shifts in the 2D as well as the G peak position, and the degree to which the peaks are shifted is distinct for each effect. In contrast, interlayer coupling modifies only the 2D peak position, allowing this effect to be distinguished from strain and doping. Despite being weak in layered materials, the van der Waals interaction is capable of modifying graphene's electronic band structure in graphene-hBN heterostructures,<sup>[43]</sup> and motivates investigation of interlayer coupling in our  $\text{MoS}_2$ -graphene heterostructures.



**Figure 5.** Raman spectroscopy of heterostructure samples. a) Comparison of Raman peaks for  $\text{MoS}_2$  synthesized on graphene (upper, red) and on  $\text{SiO}_2$  (lower, blue). The characteristic  $\text{MoS}_2$   $E_{2g}^1$  and  $A_{1g}$  peaks are identified. The peak separation is  $21.3\text{ cm}^{-1}$ , consistent with single layer  $\text{MoS}_2$  on graphene. b, c) Raman maps displaying the peak separation measured across a  $10\text{ }\mu\text{m} \times 10\text{ }\mu\text{m}$  area and a  $3\text{ mm} \times 3\text{ mm}$  area, respectively, confirming single monolayer  $\text{MoS}_2$  growth. d) Spectra of graphene peaks before (black line) and after (red line) CVD of  $\text{MoS}_2$ .

Following an analysis procedure outlined in refs [43,44], we plot the 2D vs G peak position of spectra obtained from bare graphene and  $\text{MoS}_2$ -graphene heterostructures (blue circles) in Figure 6. Bare graphene shows variations due to deviations in strain across the surface, as expected. Referring to Figure 6, if the presence of  $\text{MoS}_2$  leads to compressive strain in graphene, a shift along the direction of the black vector labeled  $\Delta_s$  is expected for the  $\text{MoS}_2$ -graphene Raman spectra. Charge doping



**Figure 6.** Analysis of Raman spectra for bare graphene and MoS<sub>2</sub>-graphene heterostructures. The position of the graphene 2D peak is plotted versus the position of the graphene G peak. Shifts along the arrows labeled  $\Delta_{i-c}$ ,  $\Delta_s$ , and  $\Delta_c$  indicate the presence of interlayer coupling, compressive strain, and hole doping, respectively. The trend observed for MoS<sub>2</sub>-graphene indicates compressive strain in the graphene following MoS<sub>2</sub> synthesis.

has a smaller affect on the shift in 2D position, and the addition of holes will produce shifts along the red vector labeled  $\Delta_c$  in Figure 6. Interlayer coupling affects only the 2D peak position, resulting in a vertical shift along the green vector labeled  $\Delta_{i-c}$  in Figure 6. As evident in Figure 6, data for both bare graphene and MoS<sub>2</sub>-graphene heterostructures are well fit by a line with a slope of 1.9 along the black vector  $\Delta_s$ . The observed behavior suggests compressive strain is the dominant factor in both bare graphene-SiO<sub>2</sub> and MoS<sub>2</sub>-graphene-SiO<sub>2</sub> heterostructures, with the effect being more pronounced with the MoS<sub>2</sub> overlayer. The additional compressive strain in MoS<sub>2</sub>-graphene films may be a result of the different coefficients of thermal expansion exhibited in graphene and MoS<sub>2</sub>. The MoS<sub>2</sub> film may expand more than graphene at the elevated synthesis temperature and gradually compresses the graphene film as the heterostructure cools to room temperature. From this analysis there is no indication that the synthesized MoS<sub>2</sub> modifies the electronic band structure of graphene.

### 3. Conclusion

In conclusion, we have demonstrated synthesis of large-area heterostructures composed of MoS<sub>2</sub> on graphene. The samples are investigated with optical microscopy, AFM, XPS, PL and Raman spectroscopy and are uniform, continuous, and atomically thin. The ability to produce large-area 2D heterostructures by direct CVD growth of MoS<sub>2</sub> on graphene provides a new avenue as opposed to the commonly used (and tedious) method of multiple mechanical exfoliations. The availability of large scale MoS<sub>2</sub>-graphene heterostructures will stimulate

investigations into photonic, electronic, and spintronic properties of this new structure and will be a benefit to the scientific community.

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